# Non-perturbative behavior of interacting Bosons at intermediate densities

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## Systems of interacting Bosons: What and why?

- Bosons are a part of the natural world:
  - ▶ Quantum particles are either Fermions or Bosons (in 3D).
  - ▶ Fermions: electrons, protons, neutrinos, etc...
  - ▶ Bosons: photons, Helium atoms, Higgs particle, etc...
- They exhibit non-trivial physical behavior at low temperatures: e.g. Bose-Einstein condensation, superfluidity, quantized vortices, etc...
- It is difficult to handle mathematically: for instance, Bose-Einstein condensation has never been proved in interacting models in the continuum at finite density.

## Repulsive Bose gas

• Potential:  $v(r) \ge 0$ ,  $\hat{v} \ge 0$  and  $v \in L_1(\mathbb{R}^3)$ , on a torus of volume V:

$$H_N := -\frac{1}{2} \sum_{i=1}^{N} \Delta_i + \sum_{1 \le i < j \le N} v(|x_i - x_j|)$$

- Smallest eigenvalue:  $E_0$  (energy), eigenvector:  $\psi_0$  (ground state) (T=0).
- Observables in the thermodynamic limit: for instance, ground state energy per particle

$$e_0 := \lim_{\substack{V,N \to \infty \\ \frac{N}{V} = \rho}} \frac{E_0}{N}.$$

• Main difficulty: dealing with the interactions.

### Known theorems

• Low density: Lee-Huang-Yang formula

$$e_0 = 2\pi\rho a \left( 1 + \frac{128}{15\sqrt{\pi}} \sqrt{\rho a^3} + o(\sqrt{\rho a^3}) \right)$$

proved: [Lieb, Yngvason, 1998], [Yau, Yin, 2009], [Fournais, Solovej, 2020], [Basti, Cenatiempo, Schlein, 2021].

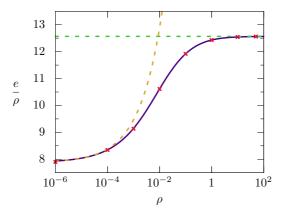
• High density: Hartree energy:

$$e_0 \sim \frac{\rho}{2} \int v$$

proved: [Lieb, 1963].

## The Simplified approach: proof of concept

For  $v(x) = e^{-|x|}$ : Simplified approach, LHY, Hartree, Monte Carlo



## The Simplified approach

- [Lieb, 1963], [Jauslin, 2023].
- Integrate  $H_N \psi_0 = E_0 \psi_0$ :

$$\int dx_1 \cdots dx_N \left( -\frac{1}{2} \sum_{i=1}^N \Delta_i \psi_0 + \sum_{1 \leqslant i < j \leqslant N} v(|x_i - x_j|) \psi_0 \right) = E_0 \int dx_1 \cdots dx_N \psi_0$$

• Therefore,

$$\frac{N(N-1)}{2} \int dx_1 dx_2 \ v(x_1 - x_2) \frac{\int dx_3 \cdots dx_N \ \psi_0}{\int dy_1 \cdots dy_N \ \psi_0} = E_0$$

## The Simplified approach

- $\psi_0 \ge 0$ , so it can be thought of as a probability distribution.
- $g_n$ : correlation functions of  $V^{-N}\psi_0$

$$g_n(x_1,\dots,x_n) := \frac{V^n \int dx_{n+1} \dots dx_N \ \psi_0(x_1,\dots,x_N)}{\int dy_1 \dots dy_N \ \psi_0(y_1,\dots,y_N)}$$

• Thus,

$$\frac{E_0}{N} = \frac{N-1}{2V} \int dx \ v(x)g_2(0,x)$$

## Hierarchy

• Equation for  $g_2$ : integrate  $H_N \psi_0 = E_0 \psi_0$  with respect to  $x_3, \dots, x_N$ :

$$-\frac{1}{2}(\Delta_x + \Delta_y)g_2(x,y) + \frac{N-2}{V} \int dz \ (v(x-z) + v(y-z))g_3(x,y,z)$$
$$+v(x-y)g_2(x,y) + \frac{(N-2)(N-3)}{2V^2} \int dz dt \ v(z-t)g_4(x,y,z,t) = E_0g_2(x,y)$$

- Infinite hierarchy of equations.
- Truncate: Factorization assumption: for n = 3, 4,

$$g_n(x_1,\dots,x_n) = \prod_{1 \le i < j \le n} (1 - u(x_i - x_j)), \quad u(x) := 1 - g_2(0,x) \in L_1(\mathbb{R}^3)$$

## Big equation

• In the thermodynamic limit,

$$-\Delta u(x) = (1 - u(x)) \left( v(x) - 2\rho K(x) + \rho^2 L(x) \right)$$

$$K := u * S, \quad S(y) := (1 - u(y))v(y)$$

$$L := u * u * S - 2u * (u(u * S)) + \frac{1}{2} \int dy dz \ u(y)u(z - x)u(z)u(y - x)S(z - y).$$

• "Big" equation:

$$L \approx u * u * S - 2u * (u(u * S))$$

## Simple equation

- Further approximate  $S(x) \approx \frac{2e}{\rho} \delta(x)$  and  $u \ll 1$ .
- Simple equation:

$$-\Delta u(x) = (1 - u(x))v(x) - 4eu(x) + 2e\rho \ u * u(x)$$
$$e = \frac{\rho}{2} \int dx \ (1 - u(x))v(x)$$

• Theorem 1: If  $v(x) \ge 0$  and  $v \in L_1 \cap L_2(\mathbb{R}^3)$ , then the simple equation has an integrable solution (proved constructively), with  $0 \le u \le 1$ .

## Energy for the simple equation

• Theorem 2:

$$\frac{e}{\rho} \xrightarrow[\rho \to \infty]{} \frac{1}{2} \int dx \ v(x).$$

This coincides with the Hartree energy.

• Theorem 3:

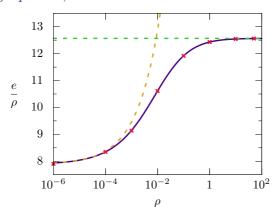
$$e = 2\pi\rho a \left( 1 + \frac{128}{15\sqrt{\pi}} \sqrt{\rho a^3} + o(\sqrt{\rho}) \right)$$

This coincides with the Lee-Huang-Yang prediction.

• **Theorem 4**: The Simple equation predicts Bose-Einstein condensation at small densities. (Agrees with Bogolyubov theory)

## Energy

 $v(x) = e^{-|x|}$ , Big equation, Monte Carlo



• Two-point correlation:

$$C_2(y-z) = \sum_{i,j} \langle \psi_0 | \delta(y-x_i) \delta(z-x_j) | \psi_0 \rangle.$$

• Radial distribution: spherical average and normalization:

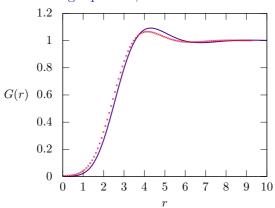
$$G(r) := \frac{1}{\rho^2} \int \frac{dy}{4\pi r^2} \, \delta(|y| - r) C_2(y).$$

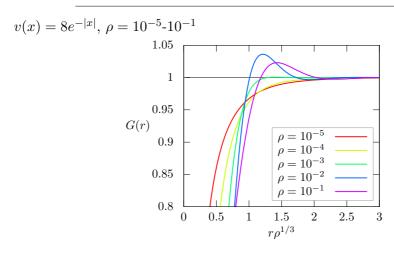
• Compute  $C_2$  using

$$C_2(x) = 2\rho \frac{\delta e_0}{\delta v(x)}.$$

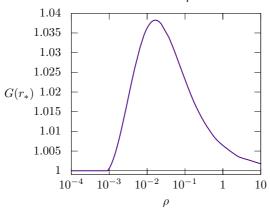
• **Theorem 5**: The Simple equation predicts that  $C_2 - 1 \sim |x|^{-4}$  for large |x|, which agrees with Bogolyubov theory.

 $v(x) = 16e^{-|x|}, \, \rho = 0.02$  Big equation, Monte Carlo





 $v(x) = 8e^{-|x|}$ , maximal value as a function of  $\rho$ :



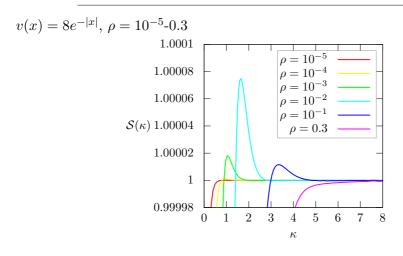
#### Structure factor

• Structure factor: Fourier transform of G:

$$S(|k|) := 1 + \rho \int dx \ e^{ikx} (G(|x|) - 1).$$

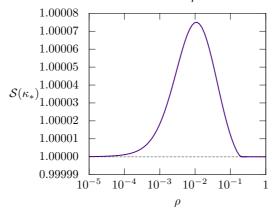
- Directly observable in X-ray scattering experiments.
- Sharp peaks: order.

### Structure factor



#### Structure factor

 $v(x) = 8e^{-|x|}$ , maximal value as a function of  $\rho$ :



## "Liquid" behavior

#### • Correlation function:

- ▶ Maximum above 1: there is a length scale at which it is more probable to find pairs of particles.
- ▶ No long range order: Short-range order.
- Structure factor:
  - ▶ Sharpening of the peak.
  - ▶ Not Bragg peaks: No long range order: Short-range order.

## Summary and outlook

- Using the Simplified approach, we were able to probe the repulsive Bose gas beyond the dilute regime.
- Evidence for non-trivial behavior at intermediate densities: short-range order.
- The intermediate density regime has not been studied much, due to the lack of tools to do so. As we have seen, there is non-trivial behavior there. This warrants further investigation, both theoretical and experimental.

## Open problems on the Simplified approach

- Connect the Simplified approach to the many-Boson system: numerics suggests the prediction of the Simplified approach is an upper bound, for all densities.
- Understand the factorization assumption. It certainly does not hold exactly. Does it hold approximately, in some sense?
- There are still many questions about the Bose gas with hard core interactions. The Simplified approach is easily defined in the hard core case. Can it shed some light?