

Non-perturbative behavior of interacting Bosons at intermediate densities

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joint with **Eric A. Carlen, Elliott H. Lieb, Markus Holzmann**

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Mathematical Physics at Rutgers Math

- **E. Carlen:** quantum mech., quantum transport, stat. mech.
- **S. Goldstein:** statistical mechanics, quantum mechanics, foundations
- **I. Jauslin:** stat mech, quantum mechanics, Renormalization Group
- **M. Kiessling:**, GR, stat. mech., quantum mech., foundations
- **J. Lebowitz:** statistical mechanics, quantum mechanics
- **A. Soffer:** scattering theory, General Relativity
- **S. Tahvildar-Zadeh:** GR, quantum mechanics, foundations
- **M. Van De Moortel:** General Relativity, Fluid dynamics

Mathematical Physics seminars

- **Mathphys seminar:** every other Thursday, 12-1,
(Elie Salomon: `ehs48@connect.rutgers.edu`)
- **Mathphys grad seminar (online):** every other Thursday, 12-1,
(Lawrence Frolov: `laf230@rutgers.edu`)
- **Mathphys webinar (online):** Wednesday, 10:45-11:45,
(Elie Salomon: `ehs48@connect.rutgers.edu`)
- **Hyperbolic PDE seminar:** Thursday, 2:30-3:30 (sometimes shifted),
(Maxime Van De Moortel: `mv715@math.rutgers.edu`)
- Graduate courses: **Intro Math Phys I** (classical mechs) & **II** (quantum mechs), **Topics in Math Phys** (GR), some years: **Statistical Mechanics**

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Systems of interacting Bosons: What and why?

- Bosons are a part of the natural world:
 - ▶ Quantum particles are either **Fermions** or **Bosons** (in 3D).
 - ▶ Fermions: electrons, protons, neutrinos, etc...
 - ▶ Bosons: photons, Helium atoms, Higgs particle, etc...
- They exhibit **non-trivial physical behavior** at low temperatures: e.g. **Bose-Einstein condensation**, superfluidity, quantized vortices, etc...
- It is **difficult to handle mathematically**: for instance, Bose-Einstein condensation has **never been proved** in interacting models in the continuum at finite density.

Repulsive Bose gas

- Potential: $v(r) \geq 0$, $\hat{v} \geq 0$ and $v \in L_1(\mathbb{R}^3)$, on a torus of volume V :

$$H_N := -\frac{1}{2} \sum_{i=1}^N \Delta_i + \sum_{1 \leq i < j \leq N} v(|x_i - x_j|)$$

- Ground state (**zero temperature**): ψ_0 , energy E_0 .
- Observables in the **thermodynamic limit**: for instance, ground state energy per particle

$$e_0 := \lim_{\substack{V,N \rightarrow \infty \\ \frac{N}{V} = \rho}} \frac{E_0}{N}.$$

- Main difficulty: dealing with the interactions.

Known theorems

- Low density: [Bogolyubov, 1947], [Lee, Huang, Yang, 1957]:

$$e_0 = 2\pi\rho a \left(1 + \frac{128}{15\sqrt{\pi}} \sqrt{\rho a^3} + o(\sqrt{\rho a^3}) \right)$$

proved: [Lieb, Yngvason, 1998], [Yau, Yin, 2009], [Fournais, Solovej, 2020], [Basti, Cenatiempo, Schlein, 2021].

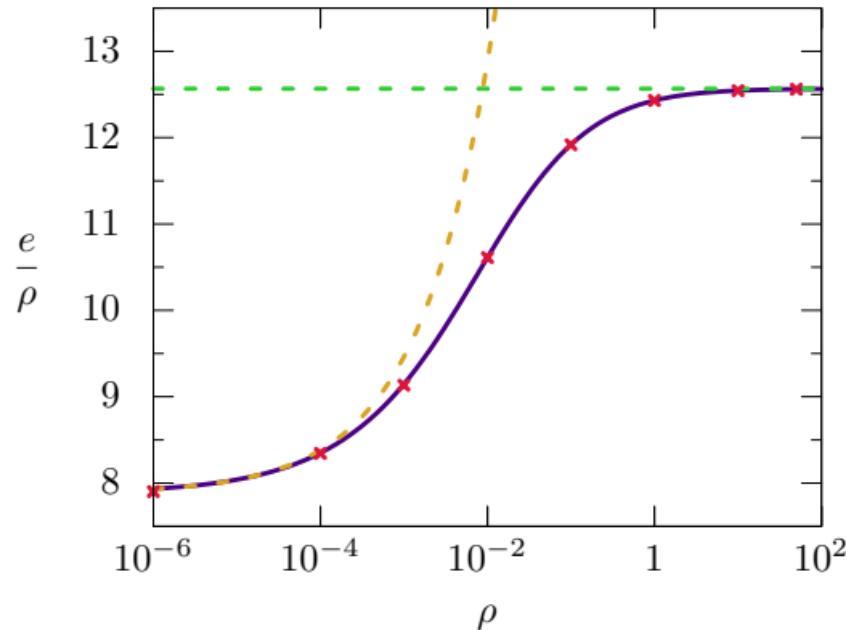
- High density: Hartree energy:

$$e_0 \sim \frac{\rho}{2} \int v$$

proved: [Lieb, 1963].

The Simplified approach: proof of concept

For $v(x) = e^{-|x|}$: Simplified approach, LHY, Hartree, Monte Carlo



The Simplified approach

- [Lieb, 1963], [Jauslin, 2023].
- Integrate $H_N \psi_0 = E_0 \psi_0$:

$$\int dx_1 \cdots dx_N \left(-\frac{1}{2} \sum_{i=1}^N \Delta_i \psi_0 + \sum_{1 \leq i < j \leq N} v(|x_i - x_j|) \psi_0 \right) = E_0 \int dx_1 \cdots dx_N \psi_0$$

- Therefore,

$$\frac{N(N-1)}{2} \int dx_1 dx_2 v(x_1 - x_2) \frac{\int dx_3 \cdots dx_N \psi_0}{\int dy_1 \cdots dy_N \psi_0} = E_0$$

The Simplified approach

- $\psi_0 \geq 0$, so it can be thought of as a probability distribution.
- g_n : correlation functions of $V^{-N}\psi_0$

$$g_n(x_1, \dots, x_n) := \frac{V^n \int dx_{n+1} \cdots dx_N \psi_0(x_1, \dots, x_N)}{\int dy_1 \cdots dy_N \psi_0(y_1, \dots, y_N)}$$

- Thus,

$$\frac{E_0}{N} = \frac{N-1}{2V} \int dx v(x) g_2(0, x)$$

Hierarchy

- Equation for g_2 : integrate $H_N\psi_0 = E_0\psi_0$ with respect to x_3, \dots, x_N :

$$\begin{aligned} & -\frac{1}{2}(\Delta_x + \Delta_y)g_2(x, y) + \frac{N-2}{V} \int dz (v(x-z) + v(y-z))g_3(x, y, z) \\ & + v(x-y)g_2(x, y) + \frac{(N-2)(N-3)}{2V^2} \int dz dt v(z-t)g_4(x, y, z, t) = E_0 g_2(x, y) \end{aligned}$$

- Infinite hierarchy of equations.

Factorization assumption

- Factorization **assumption**: for $n = 3, 4$,

$$g_n(x_1, \dots, x_n) = \prod_{1 \leq i < j \leq n} g_2(x_i - x_j), \quad u_n \in L_1(\mathbb{R}^3)$$

- In the thermodynamic limit,

$$-\Delta u(x) = (1 - u(x)) (v(x) - 2\rho K(x) + \rho^2 L(x))$$

$$K := u * S, \quad S(y) := (1 - u(y))v(y)$$

$$L := u * u * S - 2u * (u(u * S)) + \frac{1}{2} \int dy dz \, u(y)u(z-x)u(z)u(y-x)S(z-y).$$

Further simplifications

- “Big” equation:

$$L \approx u * u * S - 2u * (u(u * S))$$

- Simple equation:

$$-\Delta u(x) = (1 - u(x))v(x) - 4eu(x) + 2e\rho u * u(x)$$

$$e = \frac{\rho}{2} \int dx (1 - u(x))v(x)$$

Theorems for the simple equation

- **Theorem 1:**

$$\frac{e}{\rho} \xrightarrow{\rho \rightarrow \infty} \frac{1}{2} \int dx v(x).$$

This coincides with the Hartree energy.

- **Theorem 2:**

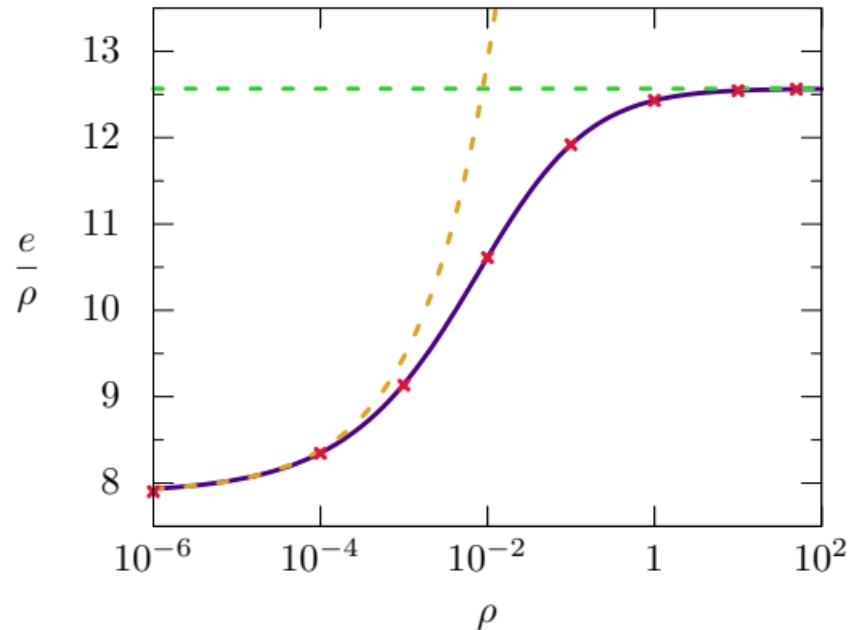
$$e = 2\pi\rho a \left(1 + \frac{128}{15\sqrt{\pi}} \sqrt{\rho a^3} + o(\sqrt{\rho}) \right)$$

This coincides with the Lee-Huang-Yang prediction.

- **Theorem 3:** The Simple equation predicts Bose-Einstein condensation at small densities.

Energy

$v(x) = e^{-|x|}$, Big equation, Monte Carlo



Radial distribution function

- Two-point correlation:

$$C_2(y - z) = \sum_{i,j} \langle \psi_0 | \delta(y - x_i) \delta(z - x_j) | \psi_0 \rangle.$$

- Radial distribution: spherical average and normalization:

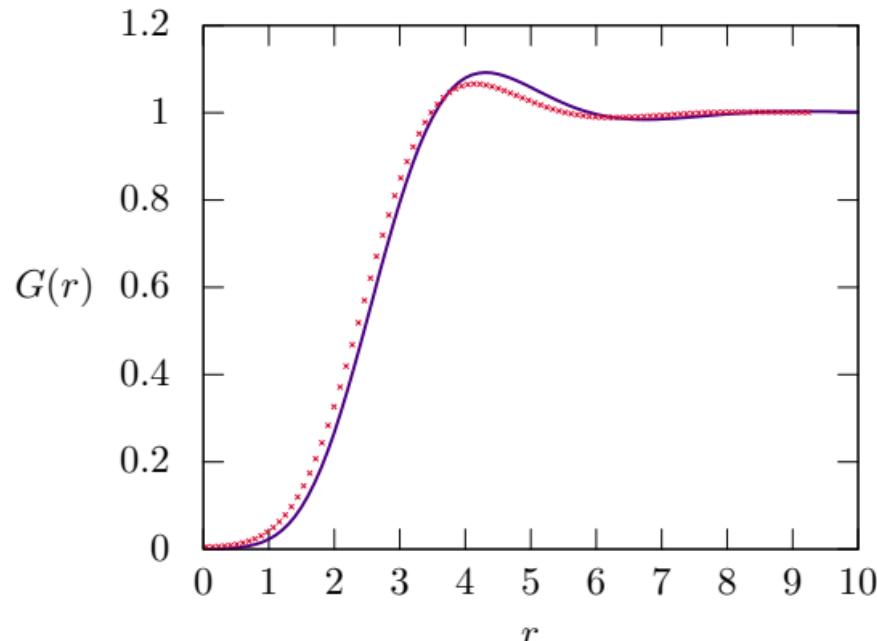
$$G(r) := \frac{1}{\rho^2} \int \frac{dy}{4\pi r^2} \delta(|y| - r) C_2(y).$$

- Compute C_2 using

$$C_2(x) = 2\rho \frac{\delta e_0}{\delta v(x)}.$$

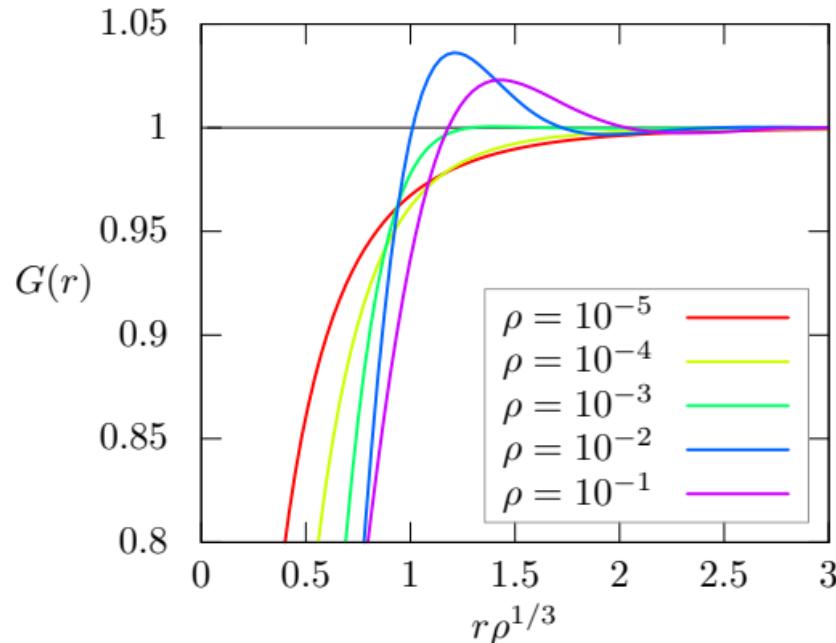
Radial distribution function

$v(x) = 16e^{-|x|}$, $\rho = 0.02$ Big equation, Monte Carlo



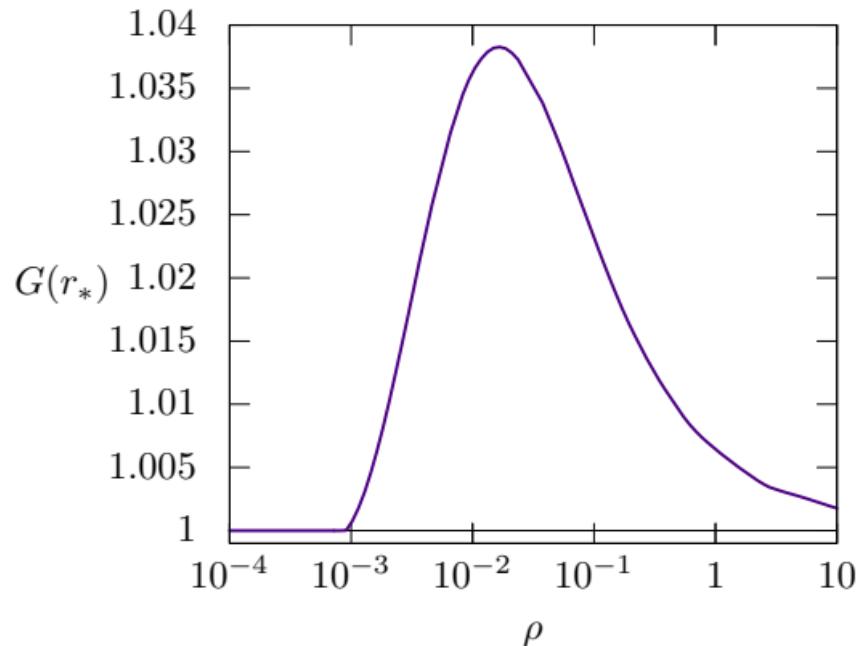
Radial distribution function

$$v(x) = 8e^{-|x|}, \rho = 10^{-5}-10^{-1}$$



Radial distribution function

$v(x) = 8e^{-|x|}$, maximal value as a function of ρ :



Liquid behavior

- Maximum above 1: there is a length scale at which it is **more probable** to find pairs of particles.
- **No** long range order.
- **Short-range order:** Liquid-like behavior.

Structure factor

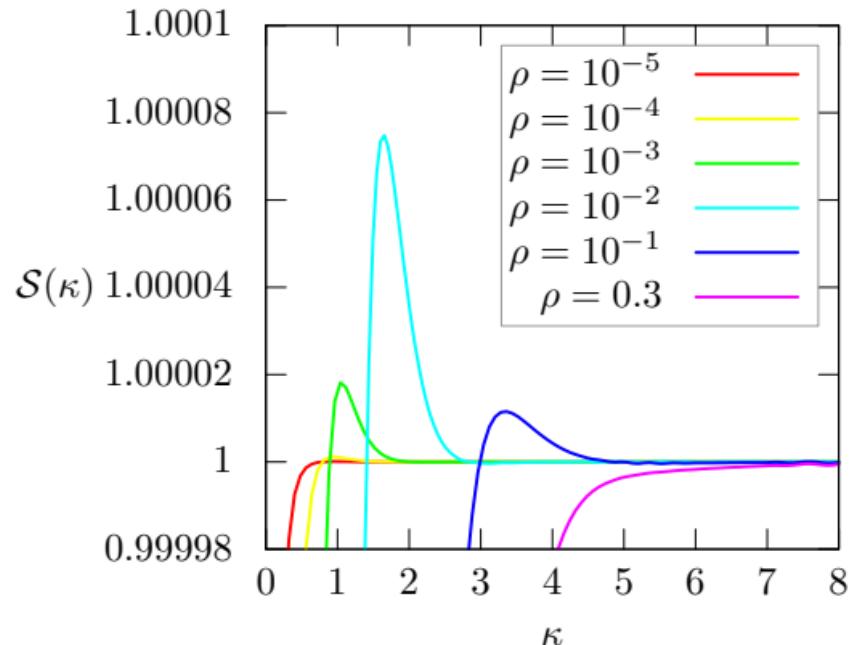
- Structure factor: Fourier transform of G :

$$S(|k|) := 1 + \rho \int dx e^{ikx} (G(|x|) - 1).$$

- Directly observable in X-ray scattering experiments.
- Sharp peaks: order.

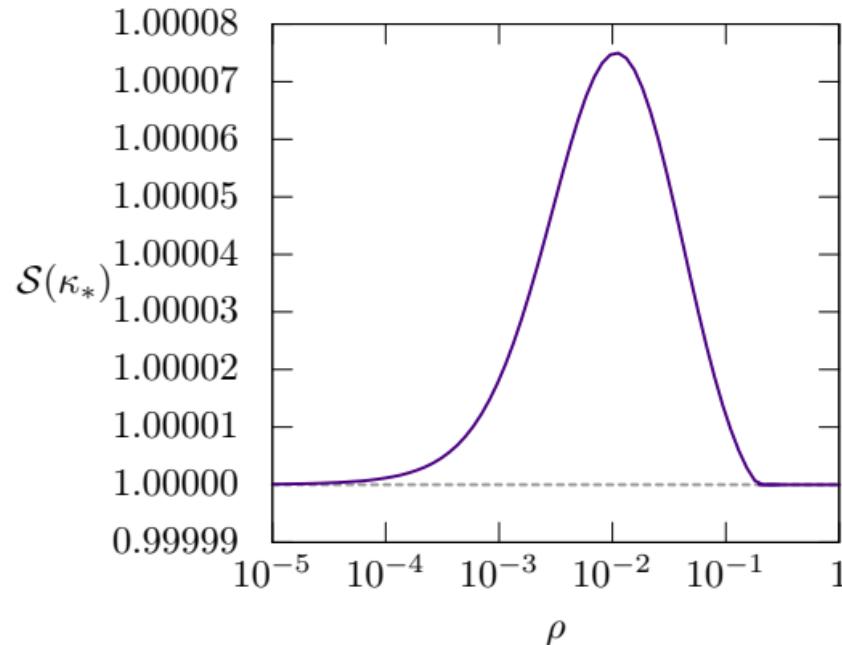
Structure factor

$$v(x) = 8e^{-|x|}, \rho = 10^{-5}-0.3$$



Structure factor

$v(x) = 8e^{-|x|}$, maximal value as a function of ρ :



Liquid behavior

- Sharpening of the peak: more order.
- Not Bragg peaks: No long range order.
- Short-range order: Liquid-like behavior.

Critical densities

- We have found two critical densities: $\rho_* \approx 0.9 \times 10^{-3}$ and $\rho_{**} \approx 0.2$.
- The radial distribution function has a maximum only for $\rho > \rho_*$.
- The structure factor has a maximum only for $\rho < \rho_{**}$.

Condensate fraction

- Proportion of particles in the condensate state:

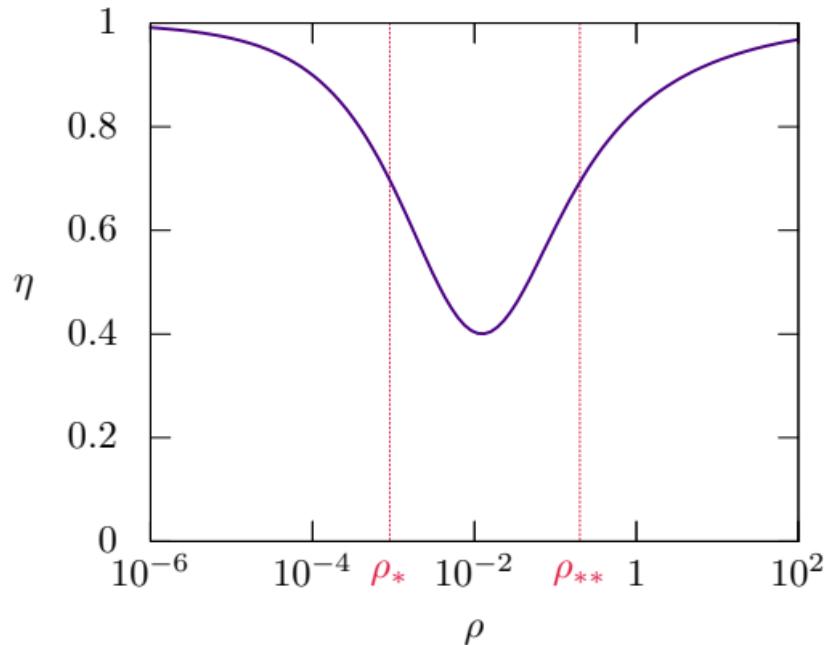
$$\eta := \frac{1}{N} \sum_i \langle \psi_0 | P_i | \psi_0 \rangle$$

where P_i is the projector onto the constant state $V^{-\frac{1}{2}}$.

- $\eta > 0$ in thermodynamic limit: **Bose-Einstein condensation** (still not proved to occur).

Condensate fraction

$$v(x) = 8e^{-|x|}:$$



Summary and outlook

- Using the Simplified approach, we were able to probe the repulsive Bose gas beyond the dilute regime.
- Evidence for non-trivial behavior at intermediate densities $\rho_* < \rho < \rho_{**}$: short-range order.
- Is there a phase transition?
- The intermediate density regime has not been studied much, due to the lack of tools to do so. As we have seen, there is non-trivial behavior there. This warrants further investigation, both theoretical and experimental.

Open problems on the Simplified approach

- Connect the Simplified approach to the many-Boson system: numerics suggests the prediction of the Simplified approach is an **upper bound**, for all densities.
- Understand the factorization assumption. It certainly does not hold exactly. Does it hold approximately, in some sense?
- There are still many questions about the Bose gas with hard core interactions. The Simplified approach is easily defined in the hard core case. Can it shed some light?